

FIRST INTEGRALS OF GEODESIC EQUATIONS  
 OF  $H$ -SPACES OF TYPE [411]

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In the present article we find 6-dimensional  $h$ -spaces of type [411] and the first integrals of geodesic equations generated by projective transformations of these spaces. The general method of finding (pseudo)Riemannian spaces which admit nonhomothetic projective group  $G_r$ , was developed by A.V. Aminova (see [1]), and is described in detail in [2] for the case of  $h$ -spaces of type [51].

For an appropriate canonical frame  $\{\xi^i\}_p$ , the Eisenhart equations for an  $h$ -space of type [411] reduce to the equation system

$$\begin{aligned} X_r \lambda_4 = 0 \text{ for } r \neq 4, \quad X_r \lambda_\sigma = 0 \text{ for } r \neq \sigma, \quad X_4(2\lambda_4 - \phi) = X_\sigma(\lambda_\sigma - 2\phi) = 0, \\ \gamma_{314} = \frac{1}{2}e_4 X_4 \phi, \quad \gamma_{413} = \gamma_{431} = \gamma_{422} = -e_4 X_4 \phi, \quad \gamma_{1\sigma 4} = \gamma_{2\sigma 3} = \gamma_{3\sigma 2} = \gamma_{4\sigma 1} = \frac{e_4 X_\sigma \phi}{\lambda_4 - \lambda_\sigma}, \\ \gamma_{2\sigma 4} = \gamma_{3\sigma 3} = \gamma_{4\sigma 2} = -\frac{e_4 X_\sigma \phi}{(\lambda_4 - \lambda_\sigma)^2}, \quad \gamma_{3\sigma 4} = \gamma_{4\sigma 3} = \frac{e_4 X_\sigma \phi}{(\lambda_4 - \lambda_\sigma)^3}, \\ \gamma_{4\sigma 4} = -\frac{e_4 X_\sigma \phi}{(\lambda_4 - \lambda_\sigma)^4}, \quad \gamma_{4\sigma \sigma} = \frac{e_\sigma X_4 \phi}{\lambda_4 - \lambda_\sigma}, \quad \gamma_{\sigma\tau\sigma} = \frac{e_\sigma X_\tau \phi}{\lambda_\sigma - \lambda_\tau}; \end{aligned} \tag{1}$$

here  $X_p \phi = \xi^i \frac{\partial \phi}{\partial x^i}$ ,  $\gamma_{pqr} = -\gamma_{qpr} = \xi_{i,j} \xi^i \xi^j$  are invariants, which play the role of the Ricci coefficients of rotation,  $\phi$  is the determining function of projective motion,  $\lambda_4, \lambda_5, \lambda_6$  are invariants such that  $\lambda_\alpha \neq \lambda_\beta$  for  $\alpha, \beta = 4, 5, 6$ ;  $e_\alpha = \pm 1$ ;  $\sigma, \tau = 5, 6$ .

The integrability conditions for system (1) can be written as follows

$$\begin{aligned} (X_1, X_2) = 0, \quad (X_1, X_3) = 0, \quad (X_2, X_3) = 0, \quad (X_1, X_4) = e_3(\gamma_{314} - \gamma_{341})X_2, \\ (X_2, X_4) = -e_2\gamma_{242}X_3, \quad (X_3, X_4) = e_3(\gamma_{134} - \gamma_{143})X_4, \quad (X_1, X_\sigma) = -e_4\gamma_{4\sigma 1}X_1, \\ (X_2, X_\sigma) = -e_3\gamma_{3\sigma 2}X_2 - e_4\gamma_{4\sigma 2}X_1, \quad (X_3, X_\sigma) = -e_2\gamma_{2\sigma 3}X_3 - e_3\gamma_{3\sigma 3}X_2 - e_4\gamma_{4\sigma 3}X_1, \\ (X_4, X_\sigma) = -e_1\gamma_{1\sigma 4}X_4 - e_2\gamma_{2\sigma 4}X_3 - e_3\gamma_{3\sigma 4}X_2 - e_4\gamma_{4\sigma 4}X_1 + e_\sigma\gamma_{\sigma 4\sigma}X_\sigma, \\ (X_\sigma, X_\tau) = -e_\sigma\gamma_{\sigma\tau\sigma}X_\sigma + e_\tau\gamma_{\tau\sigma\tau}X_\tau. \end{aligned} \tag{2}$$

By equating the coefficients at the derivatives  $\frac{\partial}{\partial x^i}$  in the left-hand and the right-hand sides of (2) we obtain a system of partial differential equations. Then, by integrating this system, we get the components  $\xi^i$  of the vectors with respect to the canonical frame. Using these results and the corresponding formulas (see [2]), we find the contravariant components of the metric tensor of the  $h$ -space under consideration, as well as the covariant components. We have

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